

EXOTIC DEFORMATION QUANTIZATION

VALENTIN OVSIENKO

1. Introduction

Let \mathcal{A} be one of the following commutative associative algebras: the algebra of all smooth functions on the plane: $\mathcal{A} = C^\infty(\mathbf{R}^2)$, or the algebra of polynomials $\mathcal{A} = \mathbf{C}[p, q]$ over \mathbf{R} or \mathbf{C} . There exists a non-trivial formal *associative* deformation of \mathcal{A} called the *Moyal \star -product* (or the standard \star -product). It is defined as an associative operation $\mathcal{A}^{\otimes 2} \rightarrow \mathcal{A}[[\hbar]]$ where \hbar is a formal variable. The explicit formula is:

$$(1) \quad F \star_{\hbar} G = FG + \sum_{k \geq 1} \frac{(i\hbar)^k}{2^k k!} \{F, G\}_k,$$

where $\{F, G\}_1 = \frac{\partial F}{\partial p} \frac{\partial G}{\partial q} - \frac{\partial F}{\partial q} \frac{\partial G}{\partial p}$ is the standard Poisson bracket, and the higher order terms are:

$$(2) \quad \{F, G\}_k = \sum_{i=0}^k (-1)^i \binom{k}{i} \frac{\partial^k F}{\partial p^{k-i} \partial q^i} \frac{\partial^k G}{\partial p^i \partial q^{k-i}}.$$

The Moyal product is the *unique* (modulo equivalence) non-trivial formal deformation of the associative algebra \mathcal{A} (see [13]).

Definition 1. A formal associative deformation of \mathcal{A} given by formula (1) is called a *\star -product* if the following hold:

1) the first order term coincides with the Poisson bracket: $\{F, G\}_1 = \{F, G\}$;

Received December 26 1995.

- 2) the higher order terms $\{F, G\}_k$ are given by differential operators vanishing on constants: $\{1, G\}_k = \{F, 1\}_k = 0$;
 3) $\{F, G\}_k = (-1)^k \{G, F\}_k$.

Definition 2. Two \star -products \star_{\hbar} and \star'_{\hbar} on \mathcal{A} are called *equivalent* if there exists a linear mapping $A_{\hbar} : \mathcal{A} \rightarrow \mathcal{A}[[\hbar]]$ such that

$$A_{\hbar}(F) = F + \sum_{k=1}^{\infty} A_k(F) \hbar^k$$

intertwining the operations \star_{\hbar} and \star'_{\hbar} : $A_{\hbar}(F) \star'_{\hbar} A_{\hbar}(G) = A_{\hbar}(F \star_{\hbar} G)$.

Consider now \mathcal{A} as a Lie algebra; the commutator is given by the Poisson bracket. The Lie algebra \mathcal{A} has a unique (modulo equivalence) non-trivial formal deformation called the *Moyal bracket* or the Moyal \star -commutator: $\{F, G\}_t = \frac{1}{\hbar}(F \star_{\hbar} G - G \star_{\hbar} F)$, where $t = -\hbar^2/2$.

The well-known De Wilde–Lecomte theorem [4] states the existence of a non-trivial \star -product for an arbitrary symplectic manifold. The theory of \star -products is a subject of *deformation quantization*. The geometrical proof of the existence theorem was given by B. Fedosov [9] (see [8] and [20] for clear explanation and survey of recent progress).

The main idea of this paper is to consider the algebra $\mathcal{F}(M)$ of functions (with singularities) on the cotangent bundle T^*M which are *Laurent polynomials* on the fibers. In contrast to the above algebra \mathcal{A} it turns out that for such algebras the standard \star -product is no more unique at least if M is one-dimensional: $\dim M = 1$.

We consider $M = S^1, \mathbf{R}$ in the real case, and $M = \mathcal{H}$ (the upper half-plane) in the holomorphic case. The main result of this paper is an explicit construction of a new \star -product on the algebra $\mathcal{F}(M)$ non-equivalent to the standard Moyal product. This \star -product is equivariant with respect to the Möbius transformations. The construction is based on the bilinear SL_2 -equivariant operations on tensor-densities on M , known as *Gordan transvectants* and *Rankin-Cohen brackets*.

We study the relations between the new \star -product and extensions of the Lie algebra $\text{Vect}(S^1)$.

The results of this paper are closely related to those of the recent work of P. Cohen, Yu. Manin and D. Zagier [3] where a one-parameter

family of associative products on the space of classical modular forms is constructed using the same SL_2 -equivariant bilinear operations.

2. Definition of the exotic \star -product

2.1 Algebras of Laurent polynomials. Let \mathcal{F} be one of the following associative algebras of functions:

$$\mathcal{F} = \mathbf{C}[p, 1/p] \otimes C^\infty(\mathbf{R}) \quad \text{or} \quad \mathcal{F} = \mathbf{C}[p, 1/p] \otimes \text{Hol}(\mathcal{H})$$

This means, it consists of functions of the type:

$$(3) \quad F(p, q) = \sum_{i=-N}^N p^i f_i(q),$$

where $f_i(q) \in C^\infty(\mathbf{R})$ in the real case, or $f_i(q)$ are holomorphic functions on the upper half-plane \mathcal{H} or $f_i \in C[q]$ (respectively).

We will also consider the algebra of polynomials: $\mathbf{C}[p, 1/p, q]$ (Laurent polynomials in p).

2.2 Transvectants. Consider the following bilinear operators on functions of one variable:

$$(4) \quad J_k^{m,n}(f, g) = \sum_{i+j=k} (-1)^i \binom{k}{i} \frac{(2m-i)!(2n-j)!}{(2m-k)!(2n-k)!} f^{(i)} g^{(j)},$$

where $f = f(z), g = g(z), f^{(i)}(z) = \frac{d^i f(z)}{dz^i}$.

These operators satisfy a remarkable property: they are equivariant under Möbius (linear-fractional) transformations. Namely, suppose that the transformation $z \mapsto \frac{az+b}{cz+d}$ (with $ad - bc = 1$) acts on the arguments as follows:

$$f(z) \mapsto f\left(\frac{az+b}{cz+d}\right)(cz+d)^{2m}, \quad g(z) \mapsto g\left(\frac{az+b}{cz+d}\right)(cz+d)^{2n},$$

then $J_k^{m,n}(f, g)$ transforms as:

$$J_k^{m,n}(f, g)(z) \mapsto J_k^{m,n}(f, g)\left(\frac{az+b}{cz+d}\right)(cz+d)^{2(m+n-k)}.$$

In other words, the operations (4) are bilinear SL_2 -equivariant mappings on tensor-densities:

$$J_k^{mn} : \mathcal{F}_m \otimes \mathcal{F}_n \rightarrow \mathcal{F}_{m+n-k},$$

where \mathcal{F}_l is the space of tensor-densities of degree $-l$: $\phi = \phi(z)(dz)^{-l}$.

The operations (4) were discovered more than one hundred years ago by Gordan [11] who called them the *transvectants*. They have been rediscovered many times: in the theory of modular functions by Rankin [18] and by Cohen [2] (so-called Rankin-Cohen brackets), in differential projective geometry by Janson and Peetre [12]. The “multi-dimensional transvectants” were defined in [14] in the context of the the Virasoro algebra and symplectic and contact geometry.

2.3 Main definition. Define the following bilinear mapping $\mathcal{F}^{\otimes 2} \rightarrow \mathcal{F}[[\hbar]]$, for $F = p^m f(q), G = p^n g(q)$, where $m, n \in \mathbf{Z}$, by putting:

$$(5) \quad F \widetilde{\star}_{\hbar} G = \sum_{k=0}^{\infty} \frac{(i\hbar)^k}{2^{2k}} p^{(m+n-k)} J_k^{m,n}(f, g),$$

Note, that the first order term coincides with the Poisson bracket.

This operation will be the main subject of this paper. We call it the *exotic \star -product*.

2.4 Remark. Another one-parameter family of operations on modular forms: $f \star^{\kappa} g = \sum_{n=0}^{\infty} t_n^{\kappa}(k, l) J_n^{kl}(f, g)$, where f and g are modular forms of weight k and l respectively, and $t_n^{\kappa}(k, l)$ are very interesting and complicated coefficients, is defined in [3].

3. Main theorems

We formulate here the main results of this paper. All the proofs will be given in Sections 4-7.

3.1 Non-equivalence. The Moyal \star -product (1) defines a non-trivial formal deformation of \mathcal{F} . We will show that the formula (5) defines a \star -product non-equivalent to the standard Moyal product.

Theorem 1. *The operation (5) is associative; it defines a formal deformation of the algebra \mathcal{F} which is not equivalent to the Moyal product.*

The associativity of the product (5) is a trivial corollary of Proposition 1 below. To prove the non-equivalence, we will use the relations with extensions of the Lie algebra of vector fields on S^1 : $\text{Vect}(S^1) \subset \mathcal{F}$ (cf. Sec.5).

It is interesting to note that the constructed \star -product is equivalent to the standard Moyal product if we consider it on the algebra $C^\infty(T^*M \setminus M)$ of all smooth functions (not only Laurent polynomials on fibers); cf. Corollary 1 below.

3.2 sl_2 -equivariance. The Lie algebra $sl_2(\mathbf{R})$ has two natural embeddings into the Poisson Lie algebra on \mathbf{R}^2 : the *symplectic Lie algebra* $sp_2(\mathbf{R}) \cong sl_2(\mathbf{R})$ generated by quadratic polynomials (p^2, pq, q^2) and another one with generators: (p, pq, pq^2) which is called the *Möbius algebra*.

It is well-known that the Moyal product (1) is the unique non-trivial formal deformation of the associative algebra of functions on \mathbf{R}^2 equivariant under the action of the symplectic algebra. This means, (1) satisfies the Leibnitz property:

$$(6) \quad \{F, G \star_{\hbar} H\} = \{F, G\} \star_{\hbar} H + G \star_{\hbar} \{F, H\},$$

where F is a quadratic polynomial (note that $\{F, G\}_t = \{F, G\}$ if F is a quadratic polynomial).

Theorem 2. *The product (5) is the unique formal deformation of the associative algebra \mathcal{F} equivariant under the action of the Möbius algebra.*

The product (5) is the unique non-trivial formal deformation of \mathcal{F} satisfying (6) for F from the Möbius sl_2 algebra.

3.3 Symplectomorphism Φ . The relation between the Moyal product and the product (5) is as follows. Consider the symplectic mapping

$$(7) \quad \Phi(p, q) = \left(\frac{p^2}{2}, \frac{q}{p} \right).$$

defined on $\mathbf{R}^2 \setminus \mathbf{R}$ in the real case and on \mathcal{H} in the complex case.

Proposition 1. *The product (5) is the Φ -conjugation of the Moyal product:*

$$(8) \quad F \tilde{\star}_{\hbar} G = F \star_{\hbar}^{\Phi} G := (F \circ \Phi \star_{\hbar} G \circ \Phi) \circ \Phi^{-1}.$$

Remark. The mapping (7) (in the complex case) can be interpreted as follows. It transforms the space of *holomorphic tensor-densities* of

degree $-k$ on \mathbf{CP}^1 to the space $\mathbf{C}^k[p, q]$ of polynomials of degree k . Indeed, there exists a natural isomorphism $z^n(dz)^{-m} \mapsto p^m q^n$ (where $m \geq 2n$) and $(p^m q^n) \circ \Phi = p^{2m-n} q^n$.

3.4 Operator formalism. The Moyal product is related to the following Weil quantization procedure. Define the following differential operators:

$$(9) \quad \begin{aligned} \widehat{p} &= i\hbar \frac{\partial}{\partial q}, \\ \widehat{q} &= q \end{aligned}$$

satisfying the canonical relation: $[\widehat{p}, \widehat{q}] = i\hbar \mathbf{I}$. Associate to each polynomial $F = F(p, q)$ the differential operator $\widehat{F} = \text{Sym}F(\widehat{p}, \widehat{q})$ symmetric in \widehat{p} and \widehat{q} . The Moyal product on the algebra of polynomials coincides with the product of differential operators: $\widehat{F} \star_{\hbar} \widehat{G} = \widehat{F\widehat{G}}$.

We will show that the \star -product (5) leads to the operators:

$$(10) \quad \begin{aligned} \widehat{p}^{\Phi} &= \left(\frac{i\hbar}{2}\right)^2 \Delta, \\ \widehat{q}^{\Phi} &= \frac{1}{4i\hbar} (\Delta^{-1} \circ A + A \circ \Delta^{-1}) \end{aligned}$$

(where $\Delta = \frac{\partial^2}{\partial q^2}$ and $A = 2q \frac{\partial}{\partial q} + 1$ is the dilation operator) also satisfying the canonical relation.

Remark that \widehat{p}^{Φ} and \widehat{q}^{Φ} given by (10) on the Hilbert space $L_2(\mathbf{R})$ are not equivalent to the operators (9) since \widehat{q}^{Φ} is symmetric but not self-adjoint (see [6] on this subject).

3.5 “Symplectomorphic” deformations. Let us consider the general situation.

Proposition 2. *Given a symplectic manifold V endowed with a \star -product \star_{\hbar} and a symplectomorphism Ψ of V , if there exists a hamiltonian isotopy of Ψ to the identity, then the Ψ -conjugate product \star_{\hbar}^{Ψ} defined according to the formula (8) is equivalent to \star_{\hbar} .*

Corollary 1. *The \star -product (5) considered on the algebra of all smooth functions $C^{\infty}(T^*\mathbf{R} \setminus \mathbf{R})$ is equivalent to the Moyal product.*

4. Möbius-invariance

In this section we prove Theorem 2. We show that the operations of transvectant (4) are Φ -conjugate of the terms of the Moyal product.

4.1 Lie algebra $\text{Vect}(\mathbf{R})$ and modules of tensor-densities. Let $\text{Vect}(\mathbf{R})$ be the Lie algebra of smooth (or polynomial) vector fields on \mathbf{R} :

$$X = X(x) \frac{d}{dx}$$

with the commutator

$$\left[X(x) \frac{d}{dx}, Y(x) \frac{d}{dx} \right] = (X(x)Y'(x) - X'(x)Y(x)) \frac{d}{dx}.$$

The natural embedding of the Lie algebra $sl_2 \subset \text{Vect}(\mathbf{R})$ is generated by the vector fields $d/dx, xd/dx, x^2d/dx$.

Define a 1-parameter family of $\text{Vect}(\mathbf{R})$ -actions on $C^\infty(\mathbf{R})$ given by

$$(11) \quad L_X^{(\lambda)} f = X(x)f'(x) - \lambda X'(x)f(x),$$

where $\lambda \in \mathbf{R}$. Geometrically, $L_X^{(\lambda)}$ is the operator of Lie derivative on *tensor-densities* of degree $-\lambda$:

$$f = f(x)(dx)^{-\lambda}.$$

Denote \mathcal{F}_λ the $\text{Vect}(\mathbf{R})$ -module structure on $C^\infty(\mathbf{R})$ given by (11).

4.2 Transvectant as a bilinear sl_2 -equivariant operator. The operations (4) can be defined as bilinear mappings on $C^\infty(\mathbf{R})$ which are sl_2 -equivariant:

Statement 4.1. For each $k = 0, 1, 2, \dots$ there exists a unique (up to a constant) bilinear sl_2 -equivariant mapping

$$\mathcal{F}_\mu \otimes \mathcal{F}_\nu \rightarrow \mathcal{F}_{\mu+\nu-k}.$$

It is given by $f \otimes g \mapsto J_k^{\mu,\nu}(f, g)$.

Proof. Straightforward (cf. [11], [12]).

4.3 Algebra \mathcal{F} as a module over $\text{Vect}(\mathbf{R})$. The Lie algebra $\text{Vect}(\mathbf{R})$ can be considered as a Lie subalgebra of \mathcal{F} . The embedding $\text{Vect}(\mathbf{R}) \subset \mathcal{F}$ is given by:

$$X(x) \frac{d}{dx} \mapsto pX(q).$$

The algebra \mathcal{F} is therefore, a $\text{Vect}(\mathbf{R})$ -module.

Lemma 4.2. *The algebra \mathcal{F} is decomposed to a direct sum of $\text{Vect}(\mathbf{R})$ -modules:*

$$\mathcal{F} = \bigoplus_{m \in \mathbf{Z}} \overline{\mathcal{F}_m}.$$

Proof. Consider the subspace of \mathcal{F} consisting of functions homogeneous of degree m in p : $F = p^m f(q)$. This subspace is a $\text{Vect}(\mathbf{R})$ -module isomorphic to \mathcal{F}_m . Indeed, $\{pX(q), p^m f(q)\} = p^m (Xf' - mX'f) = p^m L_X^{(m)} f$.

4.4 Projective property of the diffeomorphism Φ . The transvectants (4) coincide with the Φ -conjugate operators (2) from the Moyal product:

Proposition 4.3. *Let $F = p^m f(q), G = p^n g(q)$. Then*

$$(12) \quad \Phi^{*-1}\{\Phi^*F, \Phi^*G\}_k = \frac{k!}{2^k} p^{m+n-k} J_k^{m,n}(f, g).$$

Proof. The symplectomorphism Φ of \mathbf{R}^2 intertwines the symplectic algebra $sp_2 \equiv sl_2$ and the Möbius algebra: $\Phi^*(p, pq, pq^2) = (\frac{1}{2}p^2, \frac{1}{2}pq, \frac{1}{2}q^2)$. Therefore, the operation $\Phi^{*-1}\{\Phi^*F, \Phi^*G\}_k$ is Möbius-equivariant.

On the other hand, one has: $\Phi^*F = \frac{1}{2^m} p^{2m} f(\frac{q}{p})$ and $\Phi^*G = \frac{1}{2^n} p^{2n} g(\frac{q}{p})$. Since Φ^*F and Φ^*G are homogeneous of degree $2m$ and $2n$ (respectively), the function $\{\Phi^*F, \Phi^*G\}_k$ is also homogeneous of degree $2(m+n-k)$. Thus, the operation $\{F, G\}_k^\Phi = \Phi^{*-1}\{\Phi^*F, \Phi^*G\}_k$ defines a bilinear mapping on the space of tensor-densities $\mathcal{F}_m \otimes \mathcal{F}_n \rightarrow \mathcal{F}_{m+n-k}$ which is sl_2 -equivariant.

Statement 4.1 implies that it is proportional to $J_k^{m,n}$. One easily verifies the coefficient of proportionality for $F = p^m, G = p^n q^k$, to obtain the formula (12).

Proposition 4.3 is proven.

Remark. Proposition 4.3 was proven in [15]. We do not know whether this elementary fact has been mentioned by classics.

4.5 Proof of Theorem 2. Proposition 4.3 implies that the formula (5) is a Φ -conjugation of the Moyal product and is given by the formula (8).

Proposition 1 is proven.

It follows that (5) is a \star -product on \mathcal{F} equivariant under the action of the Möbius sl_2 algebra. Moreover, it is the unique \star -product with this property since the Moyal product is the unique \star -product equivariant under the action of the symplectic algebra.

Theorem 2 is proven.

5. Relation with extensions of the Lie algebra $\text{Vect}(S^1)$

We prove here that the \star -product (5) is not equivalent to the Moyal product.

Let $\text{Vect}(S^1)$ be the Lie algebra of vector fields on the circle. Consider the embedding $\text{Vect}(S^1) \subset \mathcal{F}$ given by functions on \mathbf{R}^2 of the type: $X = pX(q)$ where $X(q)$ is periodical: $X(q+1) = X(q)$.

5.1 An idea of the proof of Theorem 1. Consider the formal deformations of the *Lie algebra* \mathcal{F} associated to the \star -products (1) and (5). The restriction of the Moyal bracket to $\text{Vect}(S^1)$ is identically zero. We show that the restriction of the \star -commutator

$$\{\widetilde{F}, \widetilde{G}\}_t = \frac{1}{i\hbar}(F\widetilde{\star}_\hbar G - G\widetilde{\star}_\hbar F), \quad t = -\frac{\hbar^2}{2}$$

associated to the \star -product (5) defines a series of non-trivial extensions of the Lie algebra $\text{Vect}(S^1)$ by the modules $\mathcal{F}_k(S^1)$ of tensor-densities on S^1 of degree $-k$.

5.2 Extensions and the cohomology group $H^2(\text{Vect}(S^1); \mathcal{F}_\lambda)$. Recall that an *extension* of a Lie algebra by its module is defined by a 2-cocycle on it with values in this module. To define an extension of $\text{Vect}(S^1)$ by the module \mathcal{F}_λ one needs therefore a bilinear mapping $c : \text{Vect}(S^1)^{\otimes 2} \rightarrow \mathcal{F}_\lambda$ which satisfies the identity $\delta c = 0$:

$$c(X, [Y, Z]) + L_X^{(\lambda)} c(Y, Z) + (\text{cycle}_{X,Y,Z}) = 0.$$

(See [10]).

The cohomology group $H^2(\text{Vect}(S^1); \mathcal{F}_\lambda)$ were calculated in [19] (see [10]). This group is trivial for each value of λ except $\lambda = 0, -1, -2, -5, -7$. The explicit formulæ for the corresponding non-trivial cocycles are given in [17]. If $\lambda = -5, -7$, then $\dim H^2(\text{Vect}(S^1); \mathcal{F}_\lambda) = 1$, the cohomology group is generated by the unique (up to equivalence) non-trivial cocycle. We will obtain these cocycles from the \star -commutator.

5.3 Non-trivial cocycles on $\text{Vect}(S^1)$.

Consider the restriction of the \star -commutator $\widetilde{\{, \}_t}$ (corresponding to the \star -product (5)) to $\text{Vect}(S^1) \subset \mathcal{F}$: let

$$X = pX(q), \quad Y = pY(q);$$

then from (5) we have

$$\{X, Y\}_t = \{X, Y\} + \sum_{k=1}^{\infty} \frac{t^k}{2^{2k+1}} \frac{1}{p^{2k-1}} J_{2k+1}^{1,1}(X, Y).$$

It follows from the Jacobi identity that the first non-zero term of the series $\widetilde{\{X, Y\}}_t$ is a 2-cocycle on $\text{Vect}(S^1)$ with values in one of the $\text{Vect}(S^1)$ -modules $\mathcal{F}_k(S^1)$.

Denote for simplicity $J_{2k+1}^{1,1}$ by J_{2k+1} .

From the general formula (4) one obtains:

Lemma 5.1. *First two terms of $\widetilde{\{X, Y\}}_t$ are identically zero: $J_3(X, Y) = 0, J_5(X, Y) = 0$, the next two terms are proportional to:*

$$(13) \quad \begin{aligned} J_7(X, Y) &= X'''Y^{(IV)} - X^{(IV)}Y''', \\ J_9(X, Y) &= 2(X'''Y^{(VI)} - X^{(VI)}Y''') \\ &\quad - 9(X^{(IV)}Y^{(V)} - X^{(V)}Y^{(IV)}). \end{aligned}$$

The transvectant J_7 defines therefore a 2-cocycle. It is a remarkable fact that the same fact is true for J_9 :

Lemma 5.2. (See [17]). *The mappings*

$$J_7 : \text{Vect}(S^1)^{\otimes 2} \rightarrow \mathcal{F}_{-5} \quad \text{and} \quad J_9 : \text{Vect}(S^1)^{\otimes 2} \rightarrow \mathcal{F}_{-7}$$

are 2-cocycles on $\text{Vect}(S^1)$ representing the unique non-trivial classes of the cohomology groups $H^2(\text{Vect}(S^1); \mathcal{F}_{-5})$ and $H^2(\text{Vect}(S^1); \mathcal{F}_{-7})$ respectively.

Proof. Let us prove that J_9 is a 2-cocycle on $\text{Vect}(S^1)$. The Jacobi identity for the bracket $\{, \}_t$ implies:

$$\{X, J_9(Y, Z)\} + J_9(X, \{Y, Z\}) + J_3(X, J_7(Y, Z)) + (cycle_{X,Y,Z}) = 0$$

for any $X = pX(q), Y = pY(q), Z = pZ(q)$. One checks that the expression $J_3(X, J_7(Y, Z))$ is proportional to $X'''(Y'''Z^{(IV)} - Y^{(IV)}Z''')$, so that

$$J_3(X, J_7(Y, Z)) + (cycle_{X,Y,Z}) = 0.$$

We obtain the following relation:

$$\{X, J_9(Y, Z)\} + J_9(X, \{Y, Z\}) + (\text{cycle}_{X, Y, Z}) = 0,$$

which means that J_9 is a 2-cocycle. Indeed, recall that for any tensor density a , $\{pX, p^m a\} = p^m L_{X d/dx}^{(m)}(a)$. Thus, the last relation coincides with the relation $\delta J_9 = 0$.

Let us now show that the cocycle J_7 on $\text{Vect}(S^1)$ is not trivial. Consider a linear differential operator $A : \text{Vect}(S^1) \rightarrow \mathcal{F}_5$, given by: $A(X(q)d/dq) = (\sum_{i=0}^K a_i X^{(i)}(q))(dq)^5$. Then $\delta A(X, Y) = L_X^{(5)}A(Y) - L_Y^{(5)}A(X) - A([X, Y])$. The higher order part of this expression has a non-zero term $(5 - K)a_K X'Y^{(K)}$ and therefore $J_7 \neq \delta A$.

In the same way one proves that the cocycle J_9 on $\text{Vect}(S^1)$ is non-trivial.

Lemma 5.2 is proven.

It follows that the \star -product (5) on the algebra \mathcal{F} is not equivalent to the Moyal product.

Theorem 1 is proven.

6. Operator representation

We are looking for an linear mapping (depending on \hbar) $F \mapsto \widehat{F}^\Phi$ of the associative algebra of Laurent polynomials $\mathcal{F} = \mathbf{C}[p, 1/p, q]$ into the algebra of formal pseudodifferential operators on \mathbf{R} such that

$$\widehat{F \star_{\hbar} G}^\Phi = \widehat{F}^\Phi \widehat{G}^\Phi.$$

Recall that the algebra of Laurent polynomials $\mathbf{C}[p, 1/p, q]$ with the Moyal product is isomorphic to the associative algebra of pseudodifferential operators on \mathbf{R} with polynomial coefficients (see [1]). This isomorphism is defined on the generators $p \mapsto \widehat{p}$, $q \mapsto \widehat{q}$ by the operators (9) and $p^{-1} \mapsto \widehat{p}^{-1}$:

$$\widehat{p}^{-1} = \frac{1}{i\hbar}(\partial/\partial q)^{-1}.$$

6.1 Definition. Put:

$$(14) \quad \widehat{F}^\Phi = \widehat{\Phi^* F}.$$

Then $\widehat{F}^\Phi \widehat{G}^\Phi = \widehat{\Phi^* F \Phi^* G} = \Phi^* F \star_{\hbar} \Phi^* G = \Phi^*(F \star_{\hbar} G) = \widehat{F \star_{\hbar} G}^\Phi$.

One obtains the formulæ (10). Indeed,

$$\widehat{p}^\Phi = \widehat{p^2/2} = \frac{(i\hbar)^2}{2} \frac{\partial^2}{\partial q^2}.$$

Since $q = \frac{1}{2} \left(\left(\frac{1}{p}\right) \star_{\hbar} p q + p q \star_{\hbar} \left(\frac{1}{p}\right) \right)$, one gets:

$$\widehat{q}^\Phi = \frac{1}{4i\hbar} (\Delta \circ A + A \circ \Delta).$$

6.2 sl_2 -equivariance. For the Möbius sl_2 algebra one has:

$$\begin{aligned} \widehat{p}^\Phi &= \frac{(i\hbar)^2}{2} \frac{\partial^2}{\partial q^2}, \\ \widehat{pq}^\Phi &= \frac{i\hbar}{4} + \frac{i\hbar}{2} q \frac{\partial}{\partial q}, \\ \widehat{pq^2}^\Phi &= \frac{q^2}{2}. \end{aligned}$$

Lemma 6.1. *The mapping $F \mapsto \widehat{F}^\Phi$ satisfies the Möbius-equivariance condition:*

$$\widehat{\{X, F\}}^\Phi = [\widehat{X}^\Phi, \widehat{F}^\Phi]$$

for $X \in sl_2$.

Proof. It follows immediately from Theorem 2. Indeed, the \star -product (5) is sl_2 -equivariant (that is, satisfying the relation: $\{X, F\}_t = \{X, F\}$ for $X \in sl_2$).

Remark. Beautiful explicit formulæ for sl_2 -equivariant mappings from the space of tensor-densities to the space of pseudodifferential operators are given in [3].

7. Hamiltonian isotopy

The simple calculations below are quite standard for the cohomological technique. We need them to prove Corollary 1 of Sec. 3.

Given a symplectomorphism Ψ of a symplectic manifold V and a formal deformation $\{ , \}_t$ of the Poisson bracket on V , we prove that if Ψ is isotopic to the identity, then the formal deformation $\{ , \}_t^\Psi$ defined by:

$$\{F, G\}_t^\Psi = \Psi^{*-1} \{ \Psi^* F, \Psi^* G \}_t$$

is equivalent to $\{ , \}_t$. The similar proof is valid in the case of \star -products.

Recall that two symplectomorphisms Ψ and Ψ' of a symplectic manifold V are *isotopic* if there exists a family of functions $H_{(s)}$ on V such that the symplectomorphism $\Psi_1 \circ \Psi_2^{-1}$ is the flow of the Hamiltonian vector field with the Hamiltonian function $H_{(s)}, 0 \leq s \leq 1$.

Let $\Psi_{(s)}$ be the the flow of a family of functions $H = H_{(s)}$. We will prove that the equivalence class of the formal deformation $\{ , \}_t^{\Psi_{(s)}}$ does not depend on s .

7.1 Equivalence of homotopic cocycles. Let us first show that the cohomology class of the cocycle $C_3^{\Psi_{(s)}}$:

$$C_3^{\Psi_{(s)}}(F, G) = \Psi_{(s)}^{*-1} C_3(\Psi_{(s)}^* F, \Psi_{(s)}^* G)$$

does not depend on s . To do this, it is sufficient to prove that the derivative $\dot{C}_3 = \frac{d}{ds} C_3^{\Psi_{(s)}}|_{s=0}$ is a coboundary. One has

$$\dot{C}_3(F, G) = C_3(\{H, F\}, G) + C_3(F, \{H, G\}) - \{H, C_3(F, G)\}.$$

The relation $\delta C_3 = 0$ implies:

$$\dot{C}_3(F, G) = \{F, C_3(G, H)\} - \{G, C_3(F, H)\} - C_3(\{F, G\}, H).$$

This means, $\frac{d}{ds} C_3^{\Psi_{(s)}}|_{s=0} = \delta B_H$, where $B_H(F) = C_3(F, H)$.

7.2 General case. Let us apply the same arguments to prove that the deformations $\{ , \}_t^{\Psi_{(s)}}$ are equivalent to each other for all values of s . For this purpose we must show that there exists a family of mappings $A_{(s)}(F) = F + \sum_{k=1}^{\infty} A_{(s)_k}(F)t^k$ such that $A_{(s)}^{-1}(\{A_{(s)}(F), A_{(s)}(G)\}_t) = \{F, G\}_t$.

It is sufficient to verify the existence of a mapping $a(F) = \sum_{k=1}^{\infty} a_k(F)t^k$ (the derivative: $a(F) = d/ds(A_{(s)}(F))|_{s=s_0}$) such that

$$\frac{d}{ds} \{F, G\}_t^{\Psi_{(s)}}|_{s=s_0} = \{a(F), G\}_t + \{F, a(G)\} - a(\{F, G\}_t)$$

Since

$$\frac{d}{ds} \{F, G\}_t^{\Psi_{(s)}}|_{s=s_0} = \{\{F, H\}, G\}_t + \{F, \{G, H\}\}_t - \{\{F, G\}_t, H\},$$

from the Jacobi identity:

$$\{\{F, H\}_t, G\}_t + \{F, \{G, H\}_t\}_t - \{\{F, G\}_t, H\}_t = 0$$

one obtains that the mapping $a(F)$ can be written in the form:

$$a(F) = \sum_{k=1}^{\infty} \frac{1}{(2k+1)!} C_{2k+1}(F, H_{(s_0)}) t^k.$$

7.3 Proof of Corollary 2. Consider the \star -product (8) given by $F \star_{\hbar}^{\Phi} G$, where $F \star_{\hbar} G$ is the Moyal product (1), and $\Phi : (p, q) \mapsto (p^2/2, q/p)$. It is defined on $\mathbf{R}^2 \setminus \mathbf{R}$.

The \star -product (8) on the algebra $C^\infty(\mathbf{R}^2 \setminus \mathbf{R})$ is equivalent to the Moyal product. Indeed, the symplectomorphism Φ is isotopic to the identity in the group of all smooth symplectomorphisms of $\mathbf{R}^2 \setminus \mathbf{R}$. The isotopy is: $\Phi_s : (p, q) \mapsto (\frac{p^{1+s}}{1+s}, \frac{q}{p^s})$, where $s \in [0, 1]$.

Recall that the \star -product $F \star_{\hbar}^{\Phi} G$ on the algebra \mathcal{F} is not equivalent to the Moyal product since it coincides with the product (5).

The family Φ_s does not preserve the algebra \mathcal{F} . Theorem 1 implies that F is not isotopic to the identity in the group of symplectomorphisms of $\mathbf{R}^2 \setminus \mathbf{R}$ preserving the algebra \mathcal{F} .

8. Discussion

8.1 Difficulties in multi-dimensional case.

There exist multi-dimensional analogues of transvectants [14] and [16].

Consider the projective space \mathbf{RP}^{2n+1} endowed with the standard contact structure (or an open domain of the complex projective space \mathbf{CP}^{2n+1}). There exists an unique bilinear differential operator of order k on tensor-densities equivariant with respect to the action of the group Sp_{2n} (see [14], [16]):

$$(15) \quad J_k^{\lambda, \mu} : \mathcal{F}_\lambda \otimes \mathcal{F}_\mu \rightarrow \mathcal{F}_{\lambda+\mu-\frac{k}{n+1}},$$

where $\mathcal{F}_\lambda = \mathcal{F}_\lambda(\mathbf{P}^{2n+1})$ is the space of tensor-densities on \mathbf{P}^{2n+1} of degree $-\lambda$:

$$f = f(x_1, \dots, x_{2n+1})(dx_1 \wedge \dots \wedge dx_{2n+1})^{-\lambda}.$$

The space of tensor-densities $\mathcal{F}_\lambda(\mathbf{RP}^{2n+1})$ is *isomorphic as a module over the group of contact diffeomorphisms* to the space of homogeneous functions on \mathbf{R}^{2n+2} , and the isomorphism is given by:

$$f \mapsto F(y_1, \dots, y_{2n+2}) = y_{2n+2}^{-\lambda(n+1)} f\left(\frac{y_1}{y_{2n+2}}, \dots, \frac{y_{2n+1}}{y_{2n+2}}\right).$$

Then the operations (15) are defined as the restrictions of the terms of the standard \star -product on \mathbf{R}^{2n+2} .

The same formula (5) defines a \star -product on the space of tensor-densities on \mathbf{CP}^{2n+1} (cf. [16]). However, there is no analogues of the symplectomorphism (7). I do not know if there exists a \star -product on the Poisson algebra $\mathbf{C}[y_{2n+2}, y_{2n+2}^{-1}] \otimes C^\infty(\mathbf{RP}^{2n+1})$ non-equivalent to the standard.

8.2 Classification problem. The classification (modulo equivalence) of \star -products on the Poisson algebra \mathcal{F} is an interesting open problem. It is related to the calculation of cohomology groups $H^2(\mathcal{F}; \mathcal{F})$ and $H^3(\mathcal{F}; \mathcal{F})$. The following result was announced in [7]: $\dim H^2(\mathcal{F}; \mathcal{F}) = 2$.

Let us formulate a conjecture in the compact case. Consider the Poisson algebra $\mathcal{F}(S^1)$ of functions on $T^*S^1 \setminus S^1$ which are Laurent polynomials on the fiber: $F(p, q) = \sum_{-N \leq i \leq N} p^i f_i(q)$ where $f_i(q+1) = f_i(q)$.

Conjecture. *Every \star -product on $\mathcal{F}(S^1)$ is equivalent to (1) or (5).*

I am grateful to Yu.I. Manin who explained to me the notion of Rankin-Cohen brackets in the theory of modular forms, for some important remarks, references and clarifying discussions. It is a pleasure for me to thank O. Ogievetsky and C. Roger for multiple stimulation collaboration, and also M. Audin, C. Duval, P. Lecomte, E. Mourre, S. Tabachnikov, P. Seibt, Ya. Soibel'man and F. Ziegler for fruitful discussions and their interest in this work.

References

- [1] M. Adler, *On a trace functional for formal pseudo-differential operators and the symplectic structure of the Korteweg-de Vries type equation*, Invent. Math. **50** (1987) 219-248.

- [2] H. Cohen, *Sums involving the values at negative integers of L functions of quadratic characters*, Math. Ann. **217** (1975) 181-194.
- [3] P. Cohen, Yu. Manin & D. Zagier, *Automorphic pseudodifferential operators*, Preprint, Max Planck Inst. für Math., Bonn, 1995.
- [4] M. De Wilde & P.B.A. Lecomte, *Existence of star-products and formal deformations of the Poisson Lie algebra of arbitrary symplectic manifold*, Lett. Math. Phys. **7** (1983) 487-496.
- [5] ———, *Formal deformation of the Poisson Lie algebra of a symplectic manifold and star-products*, Deformation Theory of Algebras and Structures and Applications, Kluwer Acad. Pub., Dordrecht, 1988, 897-960.
- [6] J. Dixmier, *Sur la Relation $i(PQ - QP) = I$* , Composito Math. **13** (1956) 263-269.
- [7] A. Dzhumadil'daev, *Cohomology of Cartan-type infinite-dimensional Lie algebras*, Preprint Munich University, 1995.
- [8] C. Emrlich & A. Weinstein, *The differential geometry of Fedosov's quantization*, Lie Theory and Geometry, in honor of B. Kostant, J.-L. Brylinski and R. Brylinski eds., Progress in Math., Birkhäuser, New York.
- [9] B. V. Fedosov, *A simple geometrical construction of deformation quantization*, J. Differential Geom. **40** (1994) 213-238.
- [10] D. B. Fuchs, *Cohomology of infinite-dimensional Lie algebras*, Consultants Bureau, New York, 1987.
- [11] P. Gordan, *Invariantentheorie*, Teubner, Leipzig, 1887.
- [12] S. Janson & J. Peetre, *A new generalization of Hankel operators (the case of higher weights)*, Math. Nachr. **132** (1987) 313-328.
- [13] A. Lichnerovicz, *Déformation d'algèbres associées à une variété symplectique. (Les \ast -produits)*, Ann. Inst. Fourier (Grenoble) **32** (1982) 157-209.
- [14] V. Yu. Ovsienko, *Contact analogues of the Virasoro algebra*, Functional Anal. Appl. **24** (1990) 54-63.
- [15] O. D. Ovsienko & V. Yu. Ovsienko, *Lie derivative of order n on a line. Tensor meaning of the Gelfand-Dickey bracket*, Adv. Soviet Math. **2** (1991).
- [16] V. Yu. Ovsienko & C. Roger, *Deformations of Poisson brackets and extensions of Lie algebras of contact vector fields*, Russian Math. Surveys **47** (1992) 135-191.
- [17] ———, *Generalizations of Virasoro group and Virasoro algebra through extensions by modules of tensor-densities on S^1* , Preprint CPT-94/P.3024, to appear in Functional Anal. Appl.
- [18] R. A. Rankin, *The construction of automorphic forms from the derivatives of a given form*, J. Indian Math. Soc. **20** (1956) 103-116.

- [19] T. Tsujishita, *On the continuous cohomology of the Lie algebra of vector fields*, Proc. Japan Acad. **A53** (1977) 134-138.
- [20] A. Weinstein, *Deformation quantization*, Séminaire Bourbaki n 789, Asterisque, 1993-94.

CENTRE NATIONAL DE LA RECHERCHE SCIENTIFIQUE,
MARSEILLE